

**Problem Set III: due Monday, February 26**

- 1)
  - a) Derive the Hasegawa-Wakatani equations. Do this like the derivation for reduced MHD, but:
    - i) neglect inductive effects, and all magnetic perturbations
    - ii) retain electron pressure in the Ohm's Law
    - iii) take electrons isothermal
    - iv) derive an equation for electron density, including parallel electron flow.
  - b) Discuss the conservation properties for this system.
  - c) Derive the quasi-linear equations for the H-W system. What do they mean?
  - d) Derive the mean vorticity and particle flux.
  - f) Relate the vorticity flux to the Reynolds stress.
  
- 2)
  - a) Starting from the H-W equations, derive the Hasegawa-Mima equation in the limit  $k_{\parallel}^2 v_{Te}^2 / \omega v \rightarrow \infty$ . What is the physics of this limit?
  - b) What quantities are conserved by the Hasegawa-Mima Equation?
  - c) What are the linear waves of the H-M system? Obtain the dispersion relation.
  - d) Recover these in the H-W system, for  $k_{\parallel}^2 v_{Th}^2 / \omega v > 1$  but not infinite. Discuss your result. How does instability occur?

- 3) This problem asks you to explore the Current Convective Instability (CCI) in a homogeneous medium and its sheared field relative, the Rippling Instability.
- a) Consider first a current carrying plasma in a straight magnetic field  $\underline{B} = B_0 \hat{z}$  - i.e. ignore the poloidal field, etc. Noting that the resistivity  $\eta$  is a function of temperature (ala' Spitzer - c.f. Kulsrud 8.7), calculate the electrostatic resistive instability growth rate, assuming  $T$  evolves according to:

$$\frac{\partial T}{\partial t} + \underline{v} \cdot \underline{\nabla} T - \chi_{\parallel} \partial_z^2 T - \chi_{\perp} \nabla_{\perp}^2 T = 0$$

and the electrostatic Ohm's Law is just

$$-\partial_z \phi = \frac{1}{\eta} \frac{d\eta}{dT} \tilde{T} J_0.$$

- b) *Thoroughly* discuss the physics of this simple instability, i.e.
- what is the free energy source?
  - what is the mechanism?
  - what are the dampings and how do they restrict the unstable spectrum?
  - how does spectral asymmetry enter?
  - what is the cell structure?
- c) Use quasilinear theory and the wave breaking limit to estimate the heat flux from the C.C.I.
- d) Now, consider the instability in a *sheared* magnetic field.
- i.) What difficulties enter the analysis?
  - ii.) Resolve the difficulty by considering coupled evolution of vorticity, Ohm's Law (in electrostatic limit but with temperature fluctuations) and electron temperature. Compute the growth rate in the limit  $\chi_{\parallel}, \chi_{\perp} \rightarrow 0$ . Compute the mode width. Discuss how asymmetry enters here. Explain why.
- e) Noting that  $\chi_{\parallel} \gg \chi_{\perp}$  (why? - see Kulsrud 8.7), estimate when parallel thermal conduction becomes an important damping effect. Can  $\chi_{\parallel}$  alone ever absolutely stabilize the rippling mode?

- f) Calculate the quasilinear heat flux and use the breaking limit to estimate its magnitude.

4) *Taylor in Flatland*

Taylor awakes one morning, and finds himself in Flatland, a 2D world. Seeking to relax, he sets about reformulating his theory for that planar universe.

- a) Write down the visco-resistive 2D MHD equations, and show that *three* quadratic quantities are conserved, as  $\eta \rightarrow 0$ ,  $\nu \rightarrow 0$ .
- b) Which of these is the most likely to constrain magnetic relaxation? Argue that
- i.) the local version of this quantity is conserved for an 'flux circle', as  $\eta \rightarrow 0$ ,
  - ii.) the global version is the most "rugged", for finite  $\eta$ .
- c) Formulate a 2D Taylor Hypothesis - i.e. that magnetic energy is minimized while the quantity you identified from b.) ii.) is conserved. What equation describes this state? Show that the solution is force-free. What quantity is constant in Flatland? Hence, what is the endstate of Taylor relaxation in 2D?
- d) Consider the possibility that  $\nu \gg \eta$  in Flatland. Derive the mean field evolution equation for mean magnetic potential. Discuss!
- e) *Optional - Extra Credit* - Describe the visit of the Terrifying Torus to Flatland. How would 2D Taylor perceive this apparition?  
 N.B. You may find it useful to consult *Flatland*, by E. Abbott.

- 5) Kulsrud; Chapter 7, Problem 4

- 6) Kulsrud; Chapter 7, Problem 2

- 7) Kulsrud; Chapter 11, Problem 1. Ignore the last paragraph.

- 8) Reformulate the Sweet–Parker Reconnection problem for weak collisionality. Assume a uniform, strong guide field  $B_0 \hat{z}$  orthogonal to the plane of reconnection. What can be said about the reconnection speed? [Note: This is an open-ended problem that asks you to synthesize the stories of the current-driven ion–acoustic instability and the resulting scattering of momentum with the S–P problem. You may find it useful to consult relevant parts of Kulsrud, Chapter 14.]
- 9) a) Derive the DNLS for weakly compressible nonlinear Alfvén packets, both by heuristic and by iterative methods. For the latter, consider the fast and slow evolution of  $\tilde{B}$  using the induction equation.
- b) Consider how the DNLS might change in long ion mean free path regimes. What types of dissipation might enter, and affect the structure of the wave packet equation? [N.B. This is open-ended.]
- c) How might one treat the case where  $\beta \rightarrow 1$  ?
- d) Derive a weak turbulence analogue of the Alfvénic packet steepening story by combining:
- the effects of Alfvén wave radiation pressure on a parallel acoustic wave
  - wave kinetics for the Alfvénic packet
- Focus on the effects of refraction and note the analysis of weak Langmuir turbulence.